

Question - Obtain an expression for the specific heat of a solid according to Einstein's theory. Discuss its limitation.  
Ans - Einstein's theory of specific heat of solid  $\Rightarrow$

Einstein assumed that the atoms in a crystal vibrate independently of each other about fixed lattice points. These vibrations are all assumed to be simple harmonic all with the same frequency. The vibrations of any one atom can be split into three independent vibrations along each of the three coordinate axes. Hence a solid containing  $N$  atoms is equivalent to  $3N$  harmonic oscillators vibrating independently of each other all with the same frequency  $\nu$ . The value of this frequency depends on the strength of the restoring force.

According to quantum mechanics a single harmonic oscillator can only exist in certain discrete energy states which can be labelled  $0, 1, 2, \dots, r$ . In the state  $r$  the energy of oscillations is given by

$$E_r = h\nu \left( r + \frac{1}{2} \right) \quad \text{--- (1)}$$

Assuming Maxwell's distribution law the number of oscillation is given by

$$N_r = N_0 e^{-h\nu \left( r + \frac{1}{2} \right) / KT} \\ = N_0 y \left( r + \frac{1}{2} \right) \quad \text{--- (2)}$$

where  $y = e^{-h\nu / KT}$

The total number of oscillations given by

$$N = N_0 + N_1 + N_2 + N_3 + \dots + N_r \\ = N_0 y^{1/2} + N_0 y^{3/2} + N_0 y^{5/2} + \dots + N_0 y \left( r + \frac{1}{2} \right) + \dots \\ = N_0 y^{1/2} (1 + y + y^2 + y^3 + \dots + y^r + \dots) \\ = N_0 y^{1/2} \left( \frac{1}{1-y} \right) = N_0 y^{1/2} \quad \text{--- (3)}$$

\* Total energy of all oscillators

$$E = N_0 h\nu \left[ \frac{1}{2} + N_1 h\nu \left( 1 + \frac{1}{2} \right) + N_2 h\nu \left( 2 + \frac{1}{2} \right) + N_3 h\nu \left( 3 + \frac{1}{2} \right) + \dots \right] \\ = N_0 h\nu \left[ \frac{1}{2} y^{1/2} + \frac{3}{2} y^{3/2} + \frac{5}{2} y^{5/2} + \dots + \left( r + \frac{1}{2} \right) y \left( r + \frac{1}{2} \right) + \dots \right] \\ = N_0 h\nu \left[ \frac{1}{2} y^{1/2} + 3y^{3/2} + 5y^{5/2} + \dots + \left( r + \frac{1}{2} \right) y \left( r + \frac{1}{2} \right) + \dots \right] \\ = N_0 h\nu \left[ \frac{1}{2} y^{1/2} + \frac{3}{2} y^{3/2} + \frac{5}{2} y^{5/2} + \dots + \left( r + \frac{1}{2} \right) y \left( r + \frac{1}{2} \right) + \dots \right] \quad \text{--- (4)}$$



Hence mean energy per oscillator

$$\begin{aligned}
 \bar{E} &= \frac{E}{N} = \frac{N_0 h \nu y^{1/2}}{2} \left( \frac{1}{1-y} + \frac{2y}{(1-y)^2} \right) \frac{1-y}{N_0 y^{1/2}} \\
 &= \frac{h\nu}{2} + \frac{h\nu y}{1-y} = \frac{h\nu}{2} + \frac{h\nu e^{-h\nu/KT}}{1 - e^{-h\nu/KT}} \quad \text{--- (5)} \\
 &= \frac{h\nu}{2} + \frac{h\nu}{e^{h\nu/KT} - 1}
 \end{aligned}$$

The first term  $\frac{h\nu}{2}$  is the zero point energy which the oscillator possesses even at the absolute zero of temp. If the solid contains  $N$  atoms it has a total of  $3N$  degrees of freedom. Hence the total energy of solid is

$$U = 3NE = 3N \left[ \frac{h\nu}{2} + \frac{h\nu}{e^{h\nu/KT} - 1} \right] \quad \text{--- (6)}$$

The frequency  $\nu$  depends on the restoring force acting on the vibrating atoms. The restoring force in turn depends on the mean interatomic distance. At constant volume we may assume that the lattice spacing is constant. Hence differentiating eq<sup>n</sup> (6) w.r. to  $T$  keeping  $\nu$  constant, we get the specific heat capacity at constant volume as

$$C_V = \left( \frac{dU}{dT} \right)_V = 3NK \left( \frac{h\nu}{KT} \right)^2 \cdot \frac{e^{-\frac{h\nu}{KT}}}{e^{\frac{h\nu}{KT}} - 1}$$

$$\text{or } C_V = 3NK \frac{x^2 e^x}{(e^x - 1)^2} \quad \text{--- (7)}$$

$$\text{where } x = \frac{h\nu}{KT} = \frac{\theta}{T} \text{ and } \theta = \frac{h\nu}{K}$$

Eq<sup>n</sup> (7) is Einstein equation for the atomic heat capacity of a solid at constant volume. ' $\theta$ ' is called Einstein's temperature. This parameter  $\theta$  is chosen empirically for each solid in such a way that the predictions eq<sup>n</sup> (7) agree with experiments.

Limitations:

As  $C_V$  depends on the temp. we have considered two cases.

Case I: High temperature limit:

If  $T \gg \theta$ , i.e.  $x \ll 1$ , then

$$e^x = 1 + x = 1 + \frac{h\nu}{kT}$$

from eq<sup>n</sup> (7)

$$C_V = 3Nk \cdot \frac{x^2(1+x)}{(1+x-1)^2} = 3Nk$$

But  $Nk = R$ , the universal gas constant

$$C_V = 3R$$

This is in agreement with Dulong and Petit law which agrees with experiment at high temp. Thus at sufficiently high temp, Einsteins result reduces to Dulong and Petit's law

(ii) Low temp. limit -

If  $T \ll \theta$ , i.e.  $x \gg 1$ , then e from eq<sup>n</sup> (7)

$$C_V = 3R \left(\frac{\theta}{T}\right)^2 e^{-\left(\frac{\theta}{T}\right)}$$

This shows that as  $T$  decreases, the specific heat decreases. It becomes zero as  $T \rightarrow 0$  as required by the third law of thermodynamics.

Einsteins formula explains the specific heat ~~curve~~ curve up to a certain point. But it fails completely at very low temp. For it predicts an exponential fall whereas decreases is according to the  $T^3$  law. A weak point about Einsteins formula is the empirical nature of the Einstein temp. The nature of this parameter can not be verified by any other physical data.